

## Alcubierre metric or warp drive metric

$$ds^2 = -dt^2 + (dx - v_s(t)f(r_s)dt)^2 + dy^2 + dz^2$$

$t$  is time

$x, y, z$  are rectangular coordinates in space

$x_s(t)$  is motion of a spaceship along the  $x - axis$

$$v_s(t) = \frac{dx_s(t)}{dt}$$

$r_s$  is the distance from the spaceship position

$$\sqrt{(x - x_s(t))^2 + y^2 + z^2}$$

$$f(r_s) = \frac{\tanh(\sigma(r_s+R)) - \tanh(\sigma(r_s-R))}{2\tanh(\sigma R)}$$

$R > 0$  radius of the spherical warp bubble

$\sigma > 0$  is the bubble thickness

Citation: Varieschi, Gabriele U. and Zily Burstein (2013).

“Conformal Gravity and the Alcubierre Warp Drive Metric” *ISRN Astronomy and Astrophysics*

## anti-de Sitter metric

$$ds^2 = \cosh^2 r dt^2 - a^2(dr^2 + \sinh^2 r \sigma)$$

$t$  time

$r$  radial

$a = \sqrt{\frac{3}{\lambda}}$  where  $\lambda$  is a cosmological constant

$\sigma = d\theta^2 + \sin^2 \theta d\phi^2$  is the standard metric on the 2-sphere  $S^2$ ,  
with  $\theta, \phi$  angular.

Citation: Juan Antonio Valiente Kroon (2016). *Conformal Methods in General Relativity*, Cambridge University Press

## de Sitter metric for special relativity

$$ds^2 = dt^2 + a^2 \cosh^2\left(\frac{t}{a}\right) \bar{h}$$

$t$  time

$a = \sqrt{\frac{3}{\lambda}}$  where  $\lambda$  is a cosmological constant

$\bar{h}$  is a metric of the 3-sphere  $S^3$

Citation: Juan Antonio Valiente Kroon (2016). *Conformal Methods in General Relativity*, Cambridge University Press

## Eddington-Finkelstein coordinates

From the Schwarzschild metric, change the time coordinate  $t$  to  $u$  and  $v$ , which are defined by

$$u = t - r - 2m \log |r - 2m|$$

$$v = t + r + 2m \log |r - 2m|$$

$m$  is a mass parameter

$r$  is radial

$$ds^2 = \frac{1}{2} \left(1 - \frac{2m}{r}\right) (dudv + dvdu) - r^2 \sigma$$

$\sigma = d\theta^2 + \sin^2 \theta d\phi^2$  is the standard metric on the 2-sphere  $S^2$ , with  $\theta, \phi$  angular.

Citation: Juan Antonio Valiente Kroon (2016). *Conformal Methods in General Relativity*, Cambridge University Press

## Friedmann-Lemaître-Robertson-Walker (FLRW) metric

$$ds^2 = -dt^2 + a^2(t)d\Omega_3^2$$

$d\Omega_3^2 = d\chi^2 + f_K^2(\chi)(d\theta^2 + \sin^2 \theta d\phi^2)$  is the metric of a 3-sphere,  
a 3-plane, or a 3-hyperboloid

$f_K = \{\sin, 1, \sinh\}$  corresponding to  $K = \{-1, 0, 1\}$

$t$  is cosmic time

$a(t)$  has the dimension of a length and is a scale factor.

Citation: Deruelle, Nathalie and Jean-Philippe Uzan (2018).  
*Relativity in Modern Physics*, Oxford Graduate Texts

# Gödel metric

$$ds^2 = a^2(dx_0^2 - dx_1^2 + dx_2^2 \frac{1}{2}e^{2x_1} - dx_3^2 + 2e^{x_1}dx_0dx_2)$$

$a > 0$  is a constant

Citation: Patrick Marquet (2021). "The Exact Gödel Metric"  
*Progress in Physics*

## Gullstrand-Painlevé or Painlevé-Gullstrand coordinates

$$ds^2 = -(1 - \frac{2m}{r})d\tau^2 + 2\sqrt{\frac{2m}{r}}drd\tau + d\vec{x}^2$$

where  $t \rightarrow \tau(t, r)$  via

$$\tau = t + \int dr \frac{\sqrt{\frac{2m}{r}}}{1 - \frac{2m}{r}} = t + 2m(2\sqrt{\frac{r}{2m}} + \ln(\frac{|\sqrt{\frac{r}{2m}} - 1|}{\sqrt{\frac{r}{2m}} + 1}))$$

$t$  time

$m$  mass

$r$  radial

$\vec{x}$  spherical

Citation: Deruelle, Nathalie and Jean-Philippe Uzan (2018).  
*Relativity in Modern Physics*, Oxford Graduate Texts

# Kerr metric

$$\begin{aligned}ds^2 = & -\left(1 - \frac{2mr}{\rho^2}\right)dT^2 + \frac{\rho^2}{r^2+a^2}\left(1 + \frac{2mr}{r^2+a^2}\right)dr^2 + \rho^2d\theta^2 \\& + \left(r^2 + a^2 + \frac{2mra^2\sin^2\theta}{\rho^2}\right)\sin^2\theta d\phi^2 \\& + 4mr\left(\frac{1}{r^2+a^2}dTdr - \frac{a\sin^2\theta}{\rho^2}dTd\phi - \frac{a\sin^2\theta}{r^2+a^2}drd\phi\right)\end{aligned}$$

$T$  time

$r$  radial

$\theta, \phi$  angular

$m$  mass

$a$  angular momentum per unit mass

$$\rho^2 = r^2 + a^2 \cos^2\theta$$

Citation: Deruelle, Nathalie and Jean-Philippe Uzan (2018).

*Relativity in Modern Physics*, Oxford Graduate Texts

## Kerr-Newman rotating charged black hole metric

$$ds^2 = -\frac{\Delta}{\rho^2}(dt - a \sin^2 \theta d\phi)^2 + \frac{\sin^2 \theta}{\rho^2}((r^2 + a^2)d\phi - adt)^2 + \frac{\rho^2}{\Delta}dr^2 + \rho^2 d\theta^2$$

*t* time

*r* radial

$\theta, \phi$  angular

*m* mass

*a* angular momentum per unit mass

$\Delta = r^2 - 2mr + a^2 + q^2$ , with *q* electric charge

$\rho^2 = r^2 + a^2 \cos^2 \theta$

Citation: Deruelle, Nathalie and Jean-Philippe Uzan (2018).  
*Relativity in Modern Physics*, Oxford Graduate Texts

# Kruskal-Szekeres coordinates

$$ds^2 = r^2(d\theta^2 + \sin^2 \theta d\phi^2) - 32m^3r^{-1}e^{-\frac{r}{2m}}dudv$$

$r$  radial

$\theta, \phi$  angular

$t$  time

$m$  mass

$$u = -\sqrt{\frac{r}{2m} - 1} e^{\frac{r}{4m}} e^{-\frac{t}{2m}}$$

$$v = -\sqrt{\frac{r}{2m} - 1} e^{\frac{r}{4m}} e^{\frac{t}{2m}}$$

Citation: Hans Stephani et al. (2003). *Exact Solutions of Einstein's Field Equations*. Cambridge University Press

## Lemaître coordinates

$$ds^2 = -d\tau^2 + \frac{1}{\frac{r}{2m}} d\rho^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

$$\tau = t + \int dr \frac{\sqrt{\frac{2m}{r}}}{1 - \frac{2m}{r}} = t + 2m(2\sqrt{\frac{r}{2m}} + \ln(\frac{|\sqrt{\frac{r}{2m}} - 1|}{\sqrt{\frac{r}{2m}} + 1}))$$

$$\rho = t + \int dr \frac{1}{\sqrt{\frac{2m}{r}}(1 - \frac{2m}{r})} = t + 2m(\sqrt{\frac{r}{2m}}(2 + \frac{r}{3m}) + \ln(\frac{|\sqrt{\frac{r}{2m}} - 1|}{\sqrt{\frac{r}{2m}} + 1}))$$

$t$  time

$r$  radial

$m$  mass

$\theta, \phi$  angular

$$r(\rho, \tau) = 2m(\frac{3}{4m}(\rho - \tau)^{\frac{2}{3}})$$

Citation: Deruelle, Nathalie and Jean-Philippe Uzan (2018).  
*Relativity in Modern Physics*, Oxford Graduate Texts

# Minkowski metric/space

$$ds^2 = dt^2 - dx^2 - dy^2 - dz^2$$

$t$  is time

$x, y, z$  are rectangular coordinates in space

# Reissner-Nordström metric

$$ds^2 = -\left(1 - \frac{2m}{r} + \frac{q^2}{r^2}\right)dt^2 + \frac{dr^2}{1 - \frac{2m}{r} + \frac{q^2}{r^2}} + r^2d\Omega_2^2$$

$t$  time

$\theta, \phi$  angular

$r$  radial

$$d\Omega_2^2 = d\theta^2 + \sin^2 \theta d\phi^2$$

$m = GM$  mass of the black hole

$q = \sqrt{GM}$  electric charge of the black hole

Citation: Deruelle, Nathalie and Jean-Philippe Uzan (2018).  
*Relativity in Modern Physics*, Oxford Graduate Texts

# Rindler coordinates

$$ds^2 = -\tilde{z}^2 d\tilde{t}^2 + dx^2 + dy^2 + d\tilde{z}^2$$

$t$  time

$x, y, z$  rectangular

$$\tilde{z} = \sqrt{z^2 - t^2}$$

$$\tilde{t} = \operatorname{arctanh} \frac{t}{z}$$

Citation: Griffiths, Jerry and Jiří Podolský (2010). *Exact Space-Times in Einstein's General Relativity*, Cambridge University Press

# Schwarzschild metric

$$ds^2 = \left(1 - \frac{2m}{r}\right)dt^2 - \left(1 - \frac{2m}{r}\right)^{-1}dr^2 - r^2\sigma$$

$m$  is a mass parameter

$r$  is radial

$t$  is time

$\sigma = d\theta^2 + \sin^2 \theta d\phi^2$  is on the 2-sphere  $S^2$ , with  $\theta, \phi$  angular.

Citation: Juan Antonio Valiente Kroon (2016). *Conformal Methods in General Relativity*, Cambridge University Press

## Taub-NUT metric

$$ds^2 = -f(r)(dt + 4l \sin^2 \frac{\theta}{2} d\phi)^2 + \frac{1}{f(r)} dr^2 + (r^2 + l^2)(d\theta^2 + \sin^2 \theta d\phi^2)$$

$t$  time

$r$  radial

$\theta, \phi$  angular

$m \geq 0$  constant

$l$  constant, known as the NUT parameter

$$f(r) = \frac{r^2 - 2mr - l^2}{r^2 + l^2}$$

Citation: Griffiths, Jerry and Jiří Podolský (2010). *Exact Space-Times in Einstein's General Relativity*, Cambridge University Press

# Weyl-Lewis-Papapetrou coordinates

$$ds^2 = e^{2\psi}(dt - \omega d\phi)^2 - e^{2(\gamma-\psi)}(d\rho^2 + dz^2) - \rho^2 e^{-2\psi} d\phi^2$$

$t$  time

$\phi$  angular

$\rho$  radial

$z$  rectangular

$\omega$  angular velocity, a function of  $\rho$

$\gamma, \psi$  functions of  $\rho$

Citation: M. Sharif (2007). "Energy-Momentum Distribution of the Weyl-Lewis-Papapetrou and the Levi-Civita Metrics"  
*Brazilian Journal of Physics*

## wormhole metric

$$ds^2 = \left(-1 + \frac{r_0}{r} - \frac{\epsilon}{r^2}\right)dt^2 + \frac{1}{1 - \frac{r_0}{r}}dr^2 + r^2d\theta^2 + r^2\sin^2\theta d\phi^2$$

$t$  is time

$r$  is radial

$\theta, \phi$  are angular

$\epsilon$  is electric charge

$r_0$  is smallest radius of throat

Citation: Amy Ksir's Maple file